

Accurate Long Rollouts of Fluid Dynamics Achieved by Symmetrically-physically Constrained Neural PDE Solvers

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Outline

(1) Introduction

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Neural PDE Integration

A general PDE:
$$\frac{\partial w}{\partial t} = \mathcal{F}(t, \boldsymbol{x}, \boldsymbol{w}, \nabla \boldsymbol{w}, \nabla^2 \boldsymbol{w}, \ldots)$$
 space and time $x \in \Omega \subset \mathbb{R}^d$
 $t \in [0, T]$

 $\boldsymbol{w}(t,x) \in \mathbb{R}^m$, initial and boundary conditions: $B[\boldsymbol{w}](t,\boldsymbol{x}) = 0, \forall x \in \partial \Omega$

An update operator:
$$oldsymbol{w}(t+\Delta t,\cdot)=\mathcal{G}[oldsymbol{w}(t,\cdot)]$$

Traditional: Numerical solver

Recent: Trained neural networks (called neural PDE integration)

Why is neural PDE integration:

Makes computing fast: coarse resolution and multiple time steps Especially important for climate and ocean systems

Challenges of neural solver :

to obtain an accurate long-stable rollout generalization outside of trained data

Our scientific question:

Could we benefit from the physical and symmetry constraints for accurate long-stable rollout?

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Rotation Equivariant CNNs



New fast versions

General E(2)-Equivariant Steerable CNNs

Jenner et al. 2022

E(n)-equivariant Steerable CNNs (escnn)

Cesa et al. 2022

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Symmetries of incompressible Navier–Stokes equations

$$rac{\partial oldsymbol{u}}{\partial t} + (oldsymbol{u} \cdot
abla) oldsymbol{u} = -rac{
abla p}{
ho} + \mu
abla^2 oldsymbol{u} + oldsymbol{f}; \qquad
abla \cdot oldsymbol{u} =$$

C-grid staggering



Symmetries of INS

$$\begin{split} \mathbf{flip} : \left\{ \begin{array}{l} S_u(-F(u),F(v)) &= -F(S_u(u,v)) \\ S_v(-F(u),F(v)) &= F(S_v(u,v)) \\ \mathbf{rotation} : \left\{ \begin{array}{l} S_u(R(v),-R(u)) &= R(S_v(u,v)) \\ S_v(R(v),-R(u)) &= R(-S_u(u,v)) \\ \end{array} \right. \\ \mathbf{flip-rotation} : \left\{ \begin{array}{l} S_u(R(F(v)),-R(-F(u))) &= R(F(S_v(u,v))) \\ S_v(R(F(v)),-R(-F(u))) &= -R(-F(S_u(u,v))) \\ \end{array} \right. \end{split}$$



We build group equivariant (p4, p4m) input layers for INS

According to the symmetries of INS, we create an input layer for p4.

$$\begin{split} y_{j,0,\cdot,\cdot}^{1} &= \sum_{i=0}^{c_{in}^{u}-1} \left(W_{j,i,\cdot,\cdot}^{u} \star u_{i,\cdot,\cdot} \right) + \sum_{i=0}^{c_{in}^{v}-1} \left(W_{j,i,\cdot,\cdot}^{v} \star v_{i,\cdot,\cdot} \right) + b_{j}, \\ y_{j,1,\cdot,\cdot}^{1} &= \sum_{i=0}^{c_{in}^{u}-1} \left(R_{\text{rot}}^{90^{\circ}}(W_{j,i,\cdot,\cdot}^{v}) \star u_{i,\cdot,\cdot} \right) + \sum_{i=0}^{c_{in}^{v}-1} \left(-R_{\text{rot}}^{90^{\circ}}(W_{j,i,\cdot,\cdot}^{u}) \star v_{i,\cdot,\cdot} \right) + b_{j}, \\ y_{j,2,\cdot,\cdot}^{1} &= \sum_{i=0}^{c_{in}^{u}-1} \left(-R_{\text{rot}}^{180^{\circ}}(W_{j,i,\cdot,\cdot}^{u}) \star u_{i,\cdot,\cdot} \right) + \sum_{i=0}^{c_{in}^{v}-1} \left(-R_{\text{rot}}^{180^{\circ}}(W_{j,i,\cdot,\cdot}^{v}) \star v_{i,\cdot,\cdot} \right) + b_{j}, \\ y_{j,3,\cdot,\cdot}^{1} &= \sum_{i=0}^{c_{in}^{u}-1} \left(-R_{\text{rot}}^{270^{\circ}}(W_{j,i,\cdot,\cdot}^{v}) \star u_{i,\cdot,\cdot} \right) + \sum_{i=0}^{c_{in}^{v}-1} \left(R_{\text{rot}}^{270^{\circ}}(W_{j,i,\cdot,\cdot}^{u}) \star v_{i,\cdot,\cdot} \right) + b_{j}. \end{split}$$

We build vector output layers on C-grid

p4
$$u_{i+0.5,j} = p_{i+1,j,0} - p_{i,j,1}$$

 $v_{i,j+0.5} = p_{i,j+1,2} - p_{i,j,3}$



$$p 4m \ u_{i+0.5,j} = p_{i+1,j,1} - p_{i,j,3} + p_{i+1,j,5} - p_{i,j,7} \\ v_{i,j+0.5} = p_{i,j+1,2} - p_{i,j,4} + p_{i,j+1,6} - p_{i,j,0}$$

Physical constraints embedded into networks

Momentum conservation

 $u^{t+1} = u^t + du - \operatorname{mean}(du)$ $v^{t+1} = v^t + dv - \operatorname{mean}(dv)$

Helmholtz decomposition $\vec{v} = \nabla q + \nabla \times \vec{a}$

Learn scalar potential

$$u^{t+1} = u^t - \frac{\partial a}{\partial y}$$
$$v^{t+1} = v^t + \frac{\partial a}{\partial x}$$

Wandel et al. 2020

The symmetry- and physics-constrained neural integrator for INS



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Shallow water equations

$$\begin{aligned} \frac{\partial \vec{u}}{\partial t} &= -C_D \frac{1}{h} \vec{u} |\vec{u}| - g \nabla \zeta + a_h \nabla^2 \vec{u} \\ \frac{\partial \zeta}{\partial t} &= -\nabla \cdot (h \vec{u}) \end{aligned}$$

Table 1: Geometric and physical constraints for SWEs

	Symmetries			
Conservation laws	þ	βn	回 う 上	
None Ø	p1 /Ø	p4 /Ø	p4m/Ø	
Mass m	p1/m	p4/m	p4m/m	

Incompressible Navier–Stokes equations

$$\begin{aligned} \frac{\partial \boldsymbol{u}}{\partial t} + (\boldsymbol{u} \cdot \nabla) \boldsymbol{u} &= -\frac{\nabla p}{\rho} + \mu \nabla^2 \boldsymbol{u} + \boldsymbol{f} \\ \nabla \cdot \boldsymbol{u} &= 0 \end{aligned}$$

Table 2:	Geometric	and	physical	constraints	for
INS			57 BI		

Conservation laws	Symmetries				
	Ð	回う	$\square n / \mathbb{N}$		
None ø	p1/Ø	p4 /Ø	p4m/Ø		
Momentum pũ	p1/p1	p4/pu	p4m/ <i>p</i> u		
Mass/moment. m +ρ <i>ū</i>	$p1/m + \rho \vec{u}$	р4/m+ <i>р</i> и́	p4m/m+ $\rho \vec{u}$		

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P4m/m outperforms other networks for SWE



P4m/m+p³ achieves a long rollout for decaying turbulence



Generalization beyond the training distributions



More generalizations for SWEs

Training is one square as initial condition.





Stability of symmetry-and physics-constrained model

Using incompressible Navier–Stokes equations for test



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we present a novel architectural component that simultaneously enforces hard constraints on symmetries and conservation laws for neural PDE integrators

we find that greater enhancements are brought about by symmetry constraints, but these do not make conservation laws redundant.

Maximally constrained networks provided the best match to reference simulations, both for individual initial conditions and on a statistical level.

Thank you for your attention



Architecture

